

SNS-Pisa

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Effective theories for high scale SUSY breaking

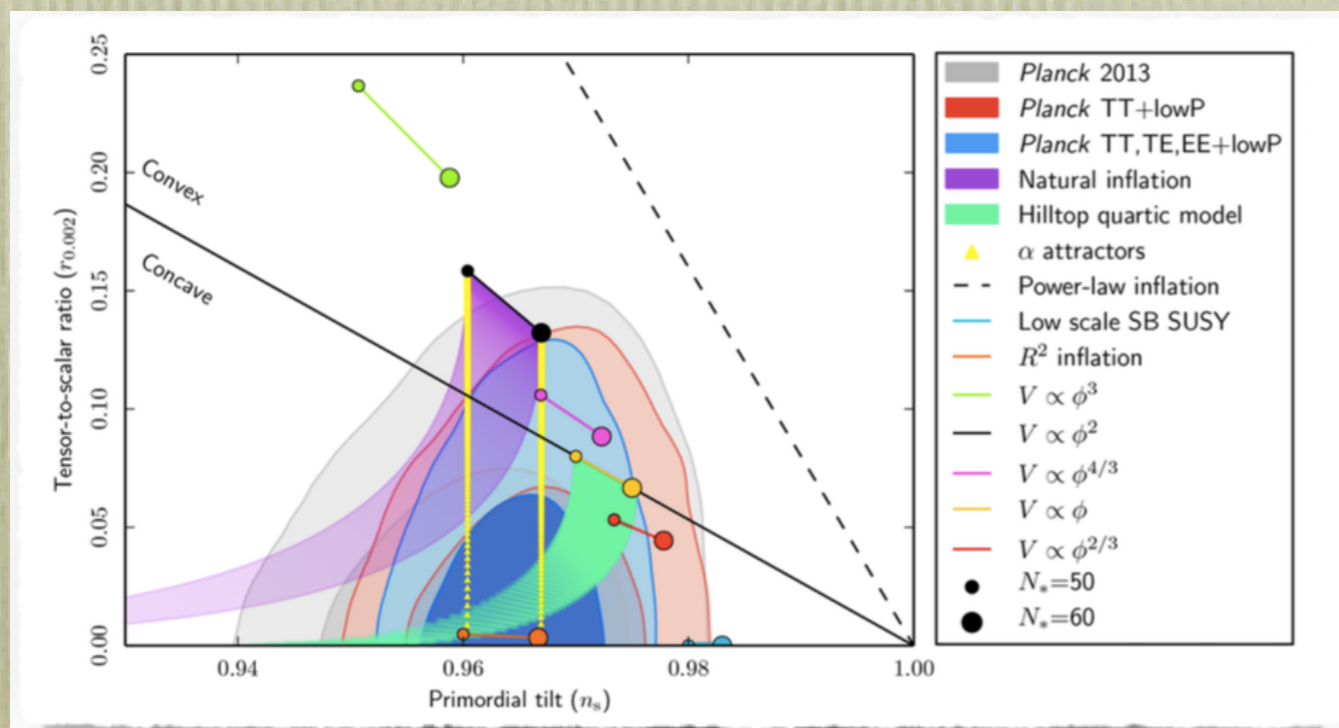
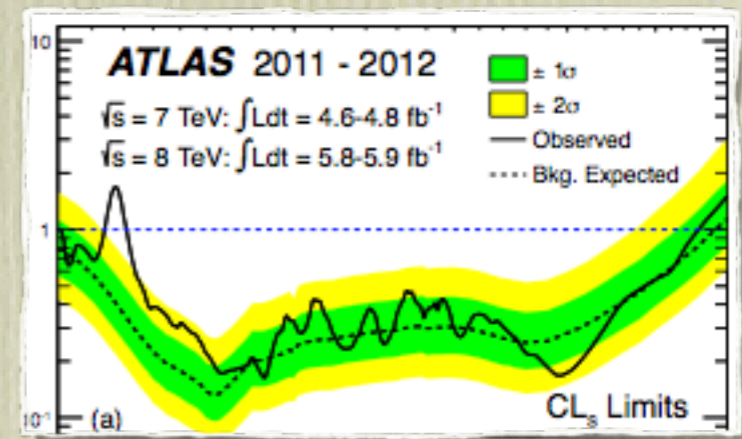
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together with:

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Standard Models: great successes, no surprises (so far)

New impressive *experimental progress* gives new *constraints* on *theoretical models* describing the origin & evolution of our Universe



Quantum Gravity necessary to UV complete inflationary models (but we may never see tensor modes)

Old & New Hierarchy problems

HUGE
(vacuum energy)

$$\langle V \rangle^{1/4} \sim 10^{-30} M_P$$

Large
(gauge hierarchy)

$$G_F^{-1/2} \sim 10^{-16} M_P$$

little

$$m_{W,Z,h}^2 \lesssim 10^{-2} m_{spart}^2$$

*LHC constrains SM-like scalar **h** at 125 GeV,
new bounds on **H, A, H[±]** and sparticles*

What about Susy?

*Supergravity relevant **iff** some superpartners (gravitino, MSSM sparticles...) are **light** with respect to KK/string/Planck cutoff scale*

*a. **Not necessarily all** superpartners*

*b. **Not necessarily at the TeV/LHC scale***

*Well motivated, but **not granted**.*

*If realized, **challenges and opportunities** for realistic supergravity models are related to BSM physics and inflation*

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- a. Not necessarily **all** superpartners*
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*If Supersymmetry is broken at **high scales** we can still use it to **constrain** and **parameterize** effective theories and provide low-energy theorems*

Non-linear symmetries

Broken symmetries useful to organize irrelevant operators using their non-linear action on the fields

*Constraints on operators of **different dimensions** connected by (non-linearly realized) symmetry*

E.g.: Higgs potential still unknown; non-linear

$SU(2) \times U(1)$ action constraints possible couplings at different order in the SMEFT expansion.

NON-LINEAR SUPERSYMMETRY

From non-linear SUSY to constrained superfields

CASALBUONI-DE
CURTIS-DOMINICI-
FERUGLIO-GATTO

Simple example (global susy): $X = x + \sqrt{2} \theta \chi + \theta^2 F^x$

$$K = |X|^2 - \frac{1}{\Lambda^2} |X|^4 \qquad W = f X$$

Linear susy in the region $\sqrt{f} \ll E \ll \Lambda$

Susy spontaneously broken at $x=0$, with susy breaking scale \sqrt{f}

Spectrum: massless goldstino + massive scalar $m = 4 \frac{f}{\Lambda}$

*At scales below m we can build an **effective Lagrangian** for the *goldstino**

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Simple example (global susy): $X = x + \sqrt{2} \theta \chi + \theta^2 F^x$

$$K = |X|^2 - \frac{1}{\Lambda^2} |X|^4 \quad W = f X$$

At zero momentum:

$$\mathcal{L} = -f^2 + |F^X + f|^2 - \frac{1}{\Lambda^2} |2x F^x - \bar{\chi} \chi|^2$$

Minimized for

$$x = \frac{\bar{\chi} \chi}{2F^x}$$

*Low-energy: **constrained chiral superfield** $X^2 = 0$*

1972 VOLKOV-AKULOV
1978 ROCEK; IVANOV-KAPUSTNIKOV;
1979 LINDSTROM-ROCEK;
1983 SAMUEL-WESS
1989 CASALBUONI-DECURTIS-DOMINICI-
FERUGLIO-GATTO;
2009 KOMARGODSKI-SEIBERG;
2011 KUZENKO-TYLER;

Chiral goldstino superfield

Various actions, all equivalent to Volkov-Akulov via non-linear field redefinitions

- *Non-linear realizations of supersymmetry:*
 - *Effective theories*
some field particularly heavy and hence integrated out
 - *Brane susy breaking models*
(high-scale induced by mutually non-susy branes)

Matter couplings using constrained superfields

*Scalarless models
(orthogonal superfields)*

$$X^2 = 0 = XY$$

BRIGNOLE-
FERUGLIO-ZWIRNER

Seiberg-Komargodski models

$$XW_\alpha = 0$$

gaugino-less

higgsino-less

$$\bar{D}_{\dot{\alpha}}(X\bar{H}) = 0$$

$$X(B - \bar{B}) = 0$$

real scalar models

Lacks organizing principle and derivation.
Supergravity realizations may change results

Supergravity realizations through Lagrange multipliers

FERRARA-KALLOSH-VAN PROEYEN

General actions with auxiliary fields

FERRARA-KALLOSH-VAN PROEYEN-WRASE
FREEDMAN-ROEST-VAN PROEYEN

Unique organizing principle and derivation from linear SUSY:

G.D., DUDAS, FARAKOS

$$X\bar{X}\Phi = 0$$

Consistently removes only the lowest component of Φ

Includes all known cases

Consistency constraints on allowed Φ

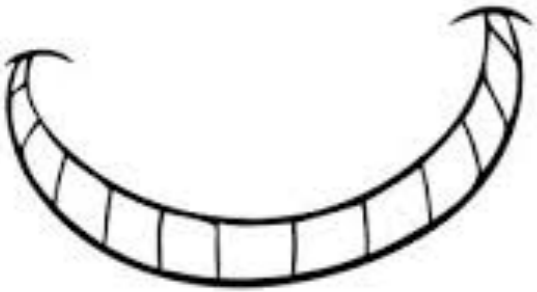
Rewriting susy lagrangian with constrained superfields

Proof that no matter the type of susy breaking (F or D term) the effective theory contains the **goldstino** constrained superfield X

Pure D-term breaking can be described by X such that $X^2 = 0$ and vector field with

$$XW_{\alpha} = 0 \quad X\bar{X}W_{\alpha} = 0$$

The latter removes the auxiliary field D, because susy breaking is now **captured by F^x**



... Sgoldstino as Cheshire cat

The constraint implies that K and W can be expanded about $X = 0$:

$$K(X, \bar{X}, z, \bar{z}) = h(z, \bar{z}) + X \bar{k}(z, \bar{z}) + \bar{X} k(z, \bar{z}) + |X|^2 g(z, \bar{z})$$

$$W(X, z) = W_0(z) + X W_1(z)$$

For the bosonic action, compute everything and set $\langle x \rangle = 0$ afterwards!

$$V = e^K (|DW|^2 - 3|W|^2)|_{x=0}$$

$$D_x W|_{x=0} = W_1(z) \neq 0$$

Nilpotent supergravity

Supergravity coupled to nilpotent field

ANTONIADIS-DUDAS-FERRARA-SAGNOTTI; BERGSHOEFF-FREEDMAN-KALLOSH-VAN PROEYEN;
HASEGAWA-YAMADA

Contains only the *graviton* and one *massive gravitino*

$$K = X \bar{X} \qquad W = m_{3/2} + f X$$

The cosmological constant is *arbitrary*

$$\Lambda = |f|^2 - 3m_{3/2}^2$$

Matter couplings

- Susy matter spectrum $m_{sp} \leq E \ll \sqrt{f}$
- *linear realizations = ordinary superfields*
- Non-susy matter spectrum $E \ll m_{sp}, \sqrt{f}$
- *non-linear realizations = Constrained superfields*

General lagrangians using previous rules

- Constrained auxiliary fields give **potentials** that *differ* from ordinary supergravity

Constrained supergravity

CRIBIORI-G.D.-FARAKOS-PORRATI

- (Old) Minimal supergravity multiplet in:

- Chiral density $2\mathcal{E} = e \{1 + i\Theta\sigma^a\bar{\psi}_a - \Theta^2(m^* + \bar{\psi}_a\bar{\sigma}^{ab}\psi_b)\}$

- Curvature superfield

$$\mathcal{R} = \left\{ m, 2\sigma^{ab}\psi_{ab} - i\sigma^a\bar{\psi}_a m + i\psi_a b^a, -\frac{1}{2}R + i\bar{\psi}^a\bar{\sigma}^b\psi_{ab} + \frac{2}{3}|m|^2 + \dots \right\}$$

- Auxiliary vector $\mathcal{D}^\alpha B_{\alpha\dot{\alpha}} = \bar{\mathcal{D}}_{\dot{\alpha}}\bar{\mathcal{R}}$.

Constrained supergravity

CRIBIORI-G.D.-FARAKOS-PORRATI

- (Old) Minimal supergravity action:

$$\mathcal{L} = \frac{3}{8} \int d^2\Theta \, 2\mathcal{E} (\overline{\mathcal{D}}^2 - 8\mathcal{R}) e^{-\frac{1}{3}K(X, \overline{X})} + \int d^2\Theta \, 2\mathcal{E} W + h.c.,$$

- Constraints on the auxiliary fields:

$$X \left(\mathcal{R} + \frac{c}{6} \right) = 0$$

Constrains auxiliary scalar

$$X \overline{X} B_{\alpha\dot{\alpha}} = 0$$

Constrains auxiliary vector

No more Kähler invariance!

Constrained supergravity

CRIBIORI-G.D.-FARAKOS-PORRATI

- Unitary gauge action:

$$e^{-1}\mathcal{L} = -\frac{1}{2}R + \frac{1}{2}\epsilon^{klmn}(\bar{\psi}_k\bar{\sigma}_l\mathcal{D}_m\psi_n - \psi_k\sigma_l\mathcal{D}_m\bar{\psi}_n) \\ - (m_{3/2}\bar{\psi}_a\bar{\sigma}^{ab}\bar{\psi}_b + \bar{m}_{3/2}\psi_a\sigma^{ab}\psi_b) - \Lambda,$$

$$\Lambda = \frac{1}{3}|c|^2 + |f|^2 + m_{3/2}\bar{c} + \bar{m}_{3/2}c = \Lambda_S - 3|m_{3/2}|^2$$

- *Susy breaking scale:* $\Lambda_S = |f|^2 + \left| \frac{c}{\sqrt{3}} + \sqrt{3}m_{3/2} \right|^2$

Matter couplings different from “de Sitter” sugra

“Anything you want” supergravity

DELACRETAZ-GORBENKO-SENATORE

CCWZ procedure to construct dressed fields

$$A = D_G [e^{\mathcal{G}Q}] \circ a$$

transforming linearly under h

$$A' = D_G [e^{\mathcal{G}'Q}] \circ a' = D_G [h(\mathcal{G}, \phi)] \circ A$$

Invariant lagrangian using ***dressed vielbein and gravitino***

No constraints on the parameters in the action:

- possible ***violation of unitarity bounds***;
- consistency of the effective action

- **Brane origin of constrained superfields**
 - *Susy is non-linearly realized on antibranes* DUDAS-MOURAD
 - *For one anti-D₃ on top of O₃ the only dof is the goldstino*
 - The anti-D₃ brane action can be written using the goldstino superfield KALLOSH-WRASE
BANDOS-MARTUCCI-SOROKIN-TONIN
BANDOS-KUZENKO-MARTUCCI-SOROKIN
 - *anti-D₃/O₃ used in **KKLT**: uplift written with manifest susy using nilpotent fields*
 - KALLOSH-LINDE
BERGSHOEFF-DASGUPTA-
KALLOSH-VANPROEYEN-WRASE
 - KALLOSH-QUEVEDO-URANGA,
APARICIO-QUEVEDO-VALANDRO,
DASGUPTA-EMELIN-MCDONOUGH
 - *More constrained multiplets* arise from *world volume fields* of the anti-D₃ brane KALLOSH-VERCNOCKE-WRASE
BANDOS-KUZENKO-MARTUCCI-SOROKIN

Summarizing:

Supersymmetry may still be relevant for pheno and cosmo if broken at "high" scale, but split spectrum

Much below heavy sparticles mass scale we should consider nonlinear realizations

*Recent developments provide **techniques**, organizing principles, **consistency conditions** and UV completions of nonlinear realizations in global and local supersymmetry*

The generation of generic EFTs with nonlinear susy from string theory and BSB is still sketchy